

Flavor Violation in SUSY GUT Scenario

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Work in Progress
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Talk Plan

- Flavor violation as a probe of new physics
- SUSY GUT Scenario and Flavor Violation
- Model-free Analysis of Flavor Violation in lepton and hadron sectors
- Numerical Results

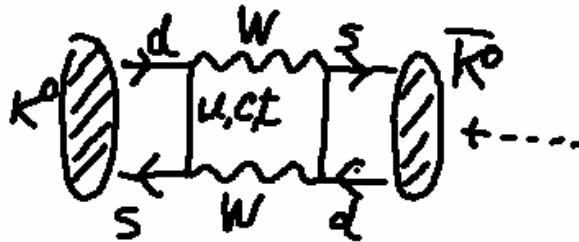
Flavor Violation as a probe of new physics

Standard Model

- Flavor Diagonal Part:
 - (spontaneously broken) gauge symmetry
- Flavor Off-Diagonal Part:
 - GIM mechanism FCNC suppression
 - loop level
 - small quark masses

Examples

$K^0 - \bar{K}^0$ mixing $\Delta m_K = 3.5 \times 10^{-12} \text{ MeV}$



$$\propto \frac{g^4}{(4\pi)^2} \frac{m_c^2}{m_W^6} \sin^2 \theta_c \cos^2 \theta_c (\bar{d}_L \gamma^\mu s_L) (\bar{d}_L \gamma_\mu s_L)$$

$\mu \rightarrow e \gamma$ $Br(\mu \rightarrow e \gamma)|_{\text{exp}} < 10^{-11}$

No lepton flavor violation in SM

One can freely rotate mass eigenbasis of massless neutrinos.

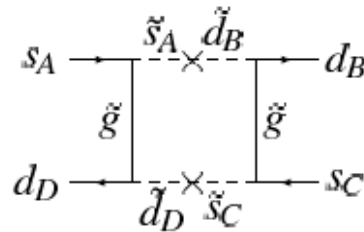
Beyond Standard Model

- New Particles and New Interactions at (sub-) TeV Scale
 - No GIM mechanism in general
- Flavor Violation:
 - good probe of new physics beyond SM

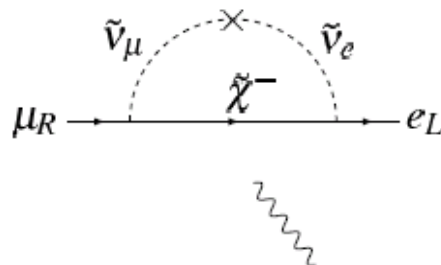
Example: Case of SUSY SM

- New particles: squarks and sleptons
- flavor mixing in squark/slepton masses
→ flavor violation

- $K^0 - \bar{K}^0$ mixing



- $\mu \rightarrow e\gamma$



Present Status of Flavor Violation

Hadron Sector:

flavor mixing (&CP violation) in CKM framework works well

K, B-mesons

Recent News: Bs-mixing is remarkable agreement with SM prediction

(Charged) Lepton Sector

No Flavor Violation observed ($\mu \rightarrow e \gamma$, $\tau \rightarrow \mu \gamma \dots$)

Only Exception:

Neutrino Oscillation: clear evidence of lepton flavor violation in neutrino sector

SUSY GUT Scenario

- Still most motivated candidate for NP beyond SM
 - low E SUSY: Sol. to gauge hierarchy problem
 - RG evolution in MSSM
 - gauge coupling unification
 - LSP: dark matter candidate

Flavor Violation in SUSY GUT Scenario

- Quarks and leptons are in the same multiplets.
- eg. SU(5)
$$10 = (Q_L, U_R^C, E_R^C)$$
$$\bar{5} = (D_R^C, L_L)$$
- “SUSY” GUT:
 - Squarks and sleptons are also in the same multiplets.
 - flavor mixing of squarks and that of sleptons are related with each other.
 - flavor violation in hadron sectors and flavor violation in lepton sectors are related.

Sources of Flavor Violation

Flavor Violation via RG evolution

- Top Yukawa & CKM above GUT scale
→ lepton flavor violation (&CPV)
- Neutrino Yukawas above GUT scale
→ hadronic flavor violation (&CPV)

Flavor Violation via Flavor Symmetry

These are sensitive to the details of GUT model

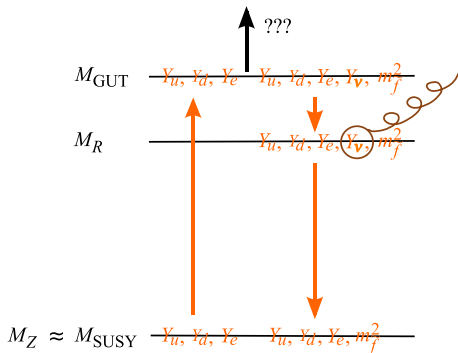
Model-free Analysis

- Generic features of SUSY GUT (SU(5))
 - At the GUT scale, squark flavor mixing is related with slepton flavor mixing.
$$d_R \leftrightarrow l_L \text{ (5-plet)}$$
$$d_L \leftrightarrow l_R \text{ (10-plet)}$$
- Questions:
 - Which gives stronger constraint on squark/slepton mass mixing, hadron sector or lepton sector?
 - Given strong constraints from hadron sector, what is the prospect of lepton flavor violation?

Numerical Results

Procedure of numerical analysis

- 1 Run Y_u, Y_d, Y_e from M_Z to M_{GUT}
- 2 Go to the basis where Y_d and Y_e are diagonal
- 3 Set Y_ν, m_f^2 at M_{GUT} **by hand**
 $A(M_{\text{GUT}}) = 0$
- 4 Run down to M_R
- 5 Remove ν_R
- 6 Run down to M_Z
- 7 Calculate μ from EWSB condition
- 8 Compute flavor violations



Mixings and constraints in consideration

- Consider $(\delta_{23}^d)_{RR} \simeq (\delta_{23}^l)_{LL}^*$ and $(\delta_{13}^d)_{LL} \simeq (\delta_{13}^l)_{RR}^*$;
for the other mixings, see e.g.

Cheung, Kang, Kim, Lee, hep-ph/0702050

Ciuchini, Masiero, Paradisi, Silvestrini, Vempati, Vives, hep-ph/0702144

- Leptonic constraints

Process	Present upper bound	Future upper bound
$B(\mu \rightarrow e\gamma)$	1.2×10^{-11}	1×10^{-13}
$B(\tau \rightarrow e\gamma)$	3.1×10^{-7}	1×10^{-8}
$B(\tau \rightarrow \mu\gamma)$	6.8×10^{-8}	1×10^{-8}

- Hadronic constraints

Observable	Range	Note
ΔM_s	$[12.5, 23] \text{ ps}^{-1}$	CDF central value $\pm 30\%$
$B(B \rightarrow X_s \gamma)$	$[2.0, 4.5] \times 10^{-4}$	
$B(B \rightarrow X_d \gamma)$	$[1, 10] \times 10^{-6}$	

Constraints on $[(\delta_{23}^d)_{RR}, (\delta_{13}^d)_{LL}]$ plane

Parameters

$$\mu > 0$$

$$\tan \beta = 3$$

$$m_0 = 210 \text{ GeV}$$

$$M_2 = 150 \text{ GeV}$$

No ν_R below M_{GUT}

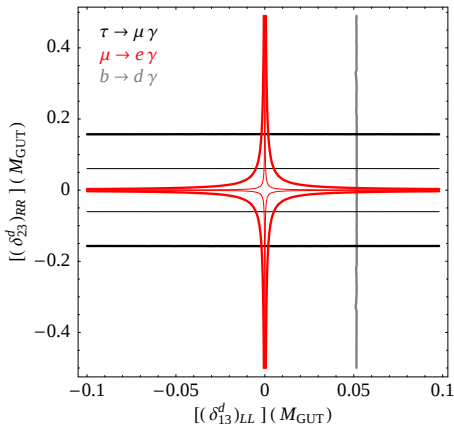
The other δ 's = 0

Constraints

$$(\delta_{13}^d)_{LL} : B \rightarrow X_d \gamma$$

$$(\delta_{23}^d)_{RR} : \Delta M_s, \tau \rightarrow \mu \gamma$$

$$(\delta_{13}^d)_{LL} \times (\delta_{23}^d)_{RR} : \mu \rightarrow e \gamma$$



Constraints on $[(\delta_{23}^d)_{RR}, (\delta_{13}^d)_{LL}]$ plane

RGE from M_* to $M_{\text{GUT}} \longrightarrow$

Parameters

$$\mu > 0$$

$$\tan \beta = 3$$

$$m_0 = 210 \text{ GeV}$$

$$M_2 = 150 \text{ GeV}$$

No ν_R below M_{GUT}

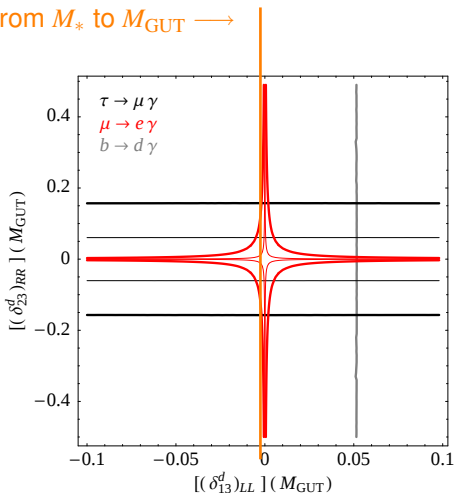
The other δ 's = 0

Constraints

$$(\delta_{13}^d)_{LL} : B \rightarrow X_d \gamma$$

$$(\delta_{23}^d)_{RR} : \Delta M_s, \tau \rightarrow \mu \gamma$$

$$(\delta_{13}^d)_{LL} \times (\delta_{23}^d)_{RR} : \mu \rightarrow e \gamma$$



Constraints on $[(\delta_{23}^d)_{RR}, (\delta_{13}^d)_{LL}]$ plane

RGE from M_* to $M_{\text{GUT}} \longrightarrow$

- Parameters

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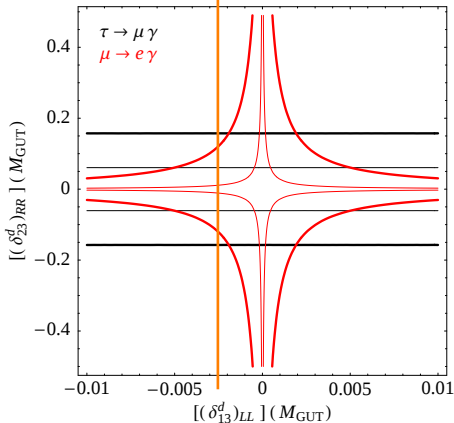
The other δ 's = 0

- Constraints

$$(\delta_{13}^d)_{LL} : B \rightarrow X_d \gamma$$

$$(\delta_{23}^d)_{RR} : \Delta M_s, \tau \rightarrow \mu \gamma$$

$$(\delta_{13}^d)_{LL} \times (\delta_{23}^d)_{RR} : \mu \rightarrow e \gamma$$



Constraints on $[\text{Re}(\delta_{23}^d)_{RR}, \text{Im}(\delta_{23}^d)_{RR}]$ plane

Parameters

$$\mu > 0$$

$$\tan\beta = 3$$

$$m_0 = 210 \text{ GeV}$$

$$M_2 = 150 \text{ GeV}$$

No ν_R below M_{GUT}

$$(\delta_{13}^d)_{LL} = 0.0025$$

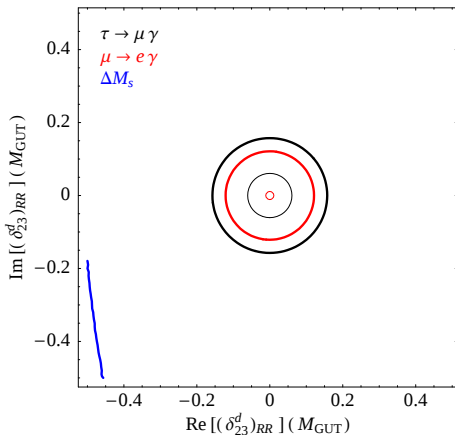
Constraints

Present

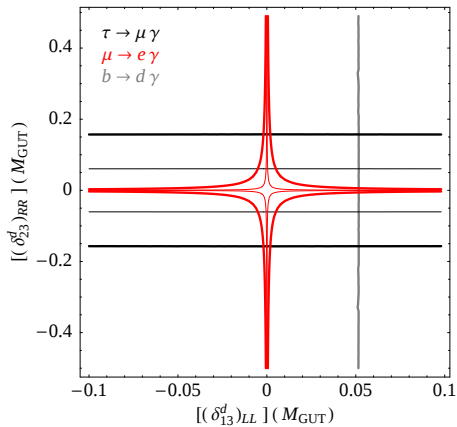
$\mu \rightarrow e\gamma$ comparable to $\tau \rightarrow \mu\gamma$

Future

$\mu \rightarrow e\gamma$ stronger than $\tau \rightarrow \mu\gamma$

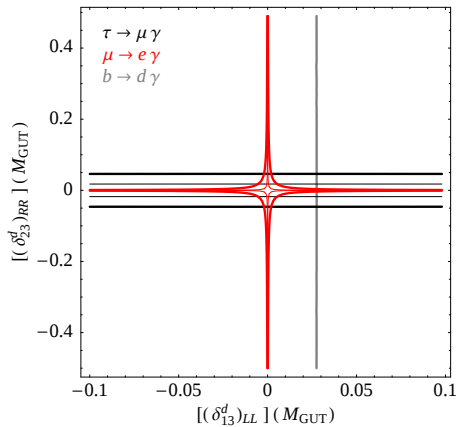


Results for higher $\tan\beta$

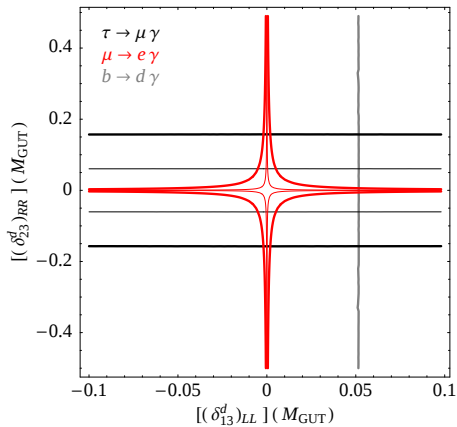


$$\tan\beta = 3$$

Results for higher $\tan \beta$

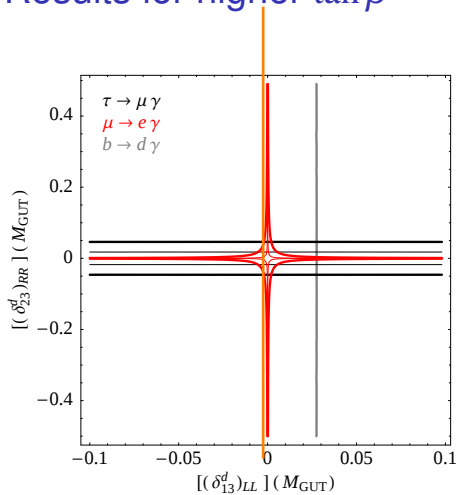


$\tan \beta = 10$

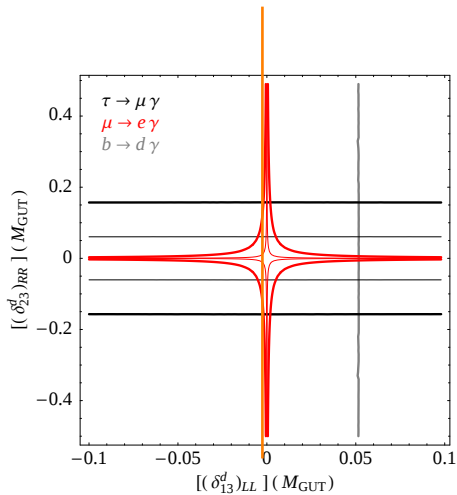


$\tan \beta = 3$

Results for higher $\tan \beta$

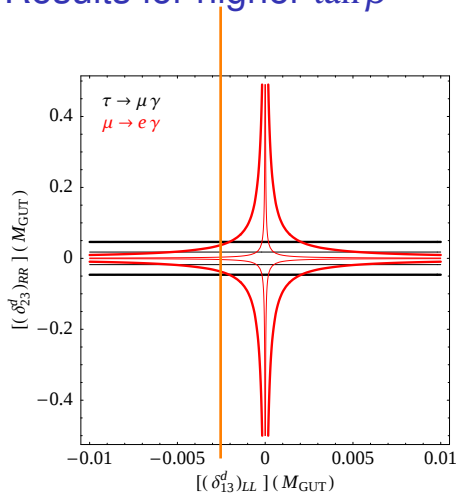


$\tan \beta = 10$

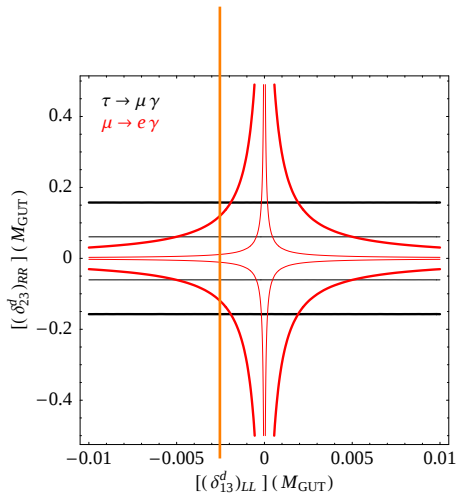


$\tan \beta = 3$

Results for higher $\tan \beta$

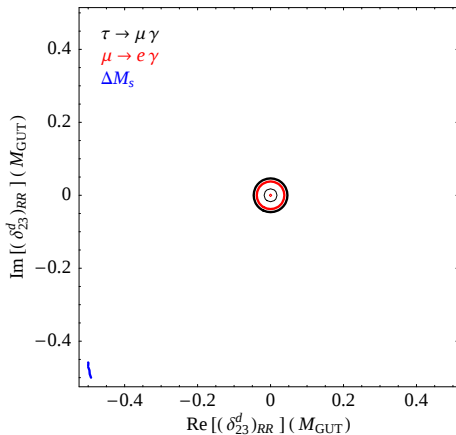


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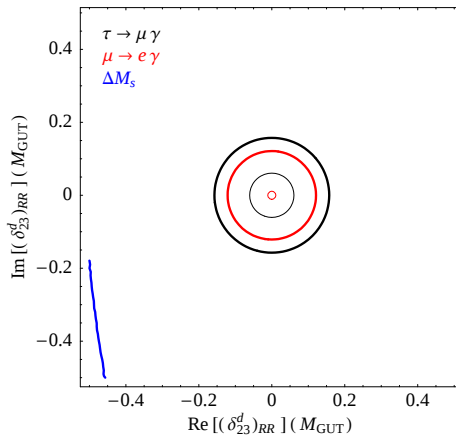


$\tan \beta = 3$

Results for higher $\tan \beta$

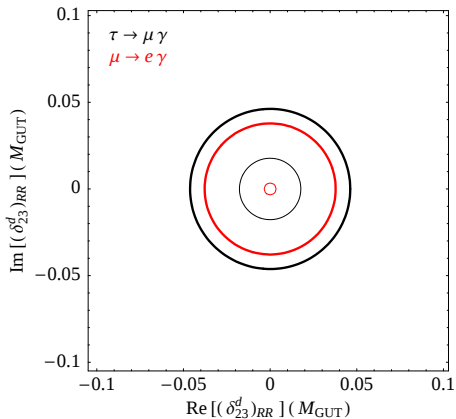


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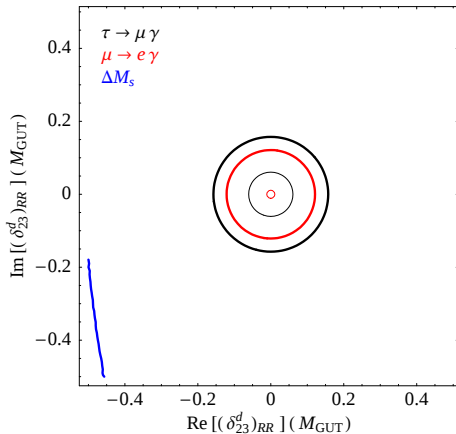


$\tan \beta = 3$

Results for higher $\tan \beta$

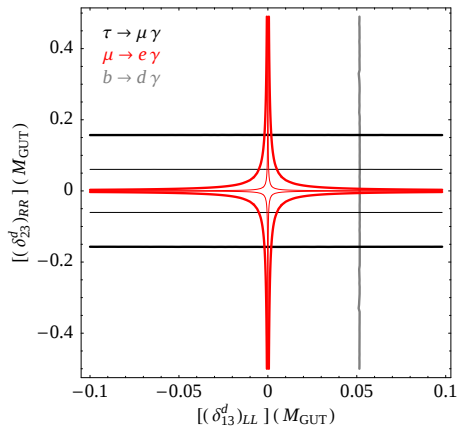


$\tan \beta = 10$



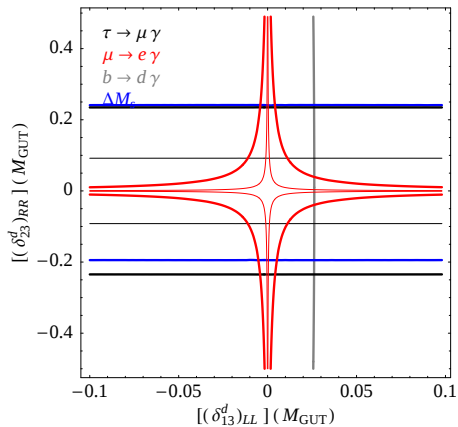
$\tan \beta = 3$

Gaugino/scalar mass dependence

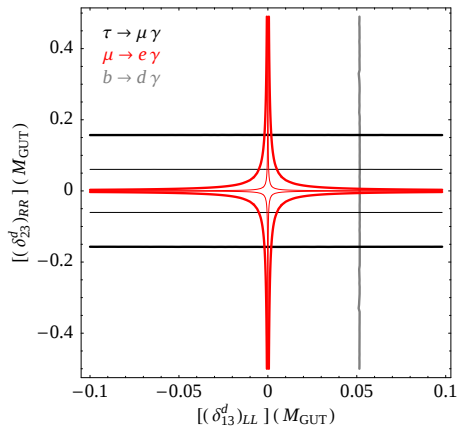


$$m_0 = 210 \text{ GeV}$$

Gaugino/scalar mass dependence

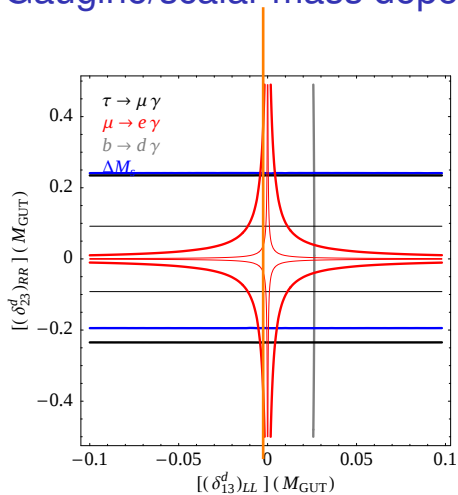


$m_0 = 400 \text{ GeV}$

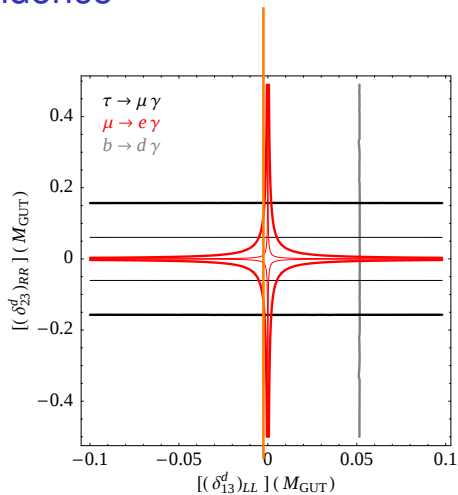


$m_0 = 210 \text{ GeV}$

Gaugino/scalar mass dependence

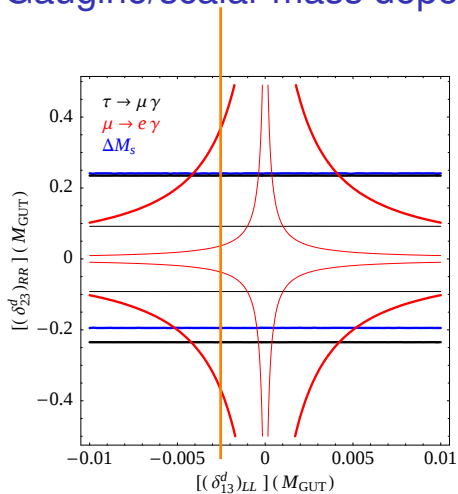


$m_0 = 400 \text{ GeV}$

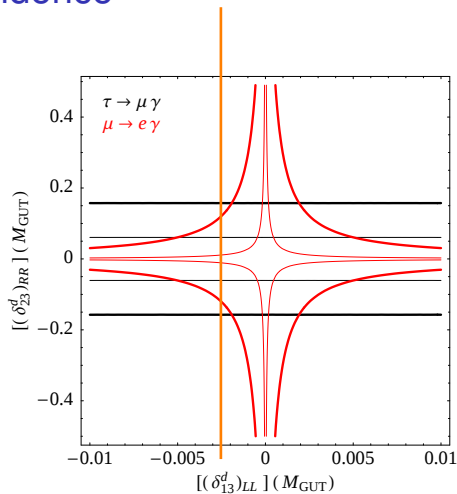


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Gaugino/scalar mass dependence

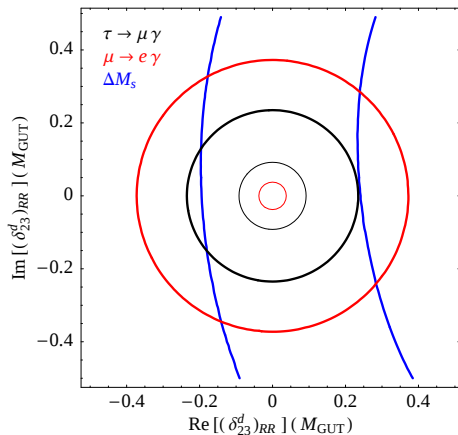


$m_0 = 400$ GeV

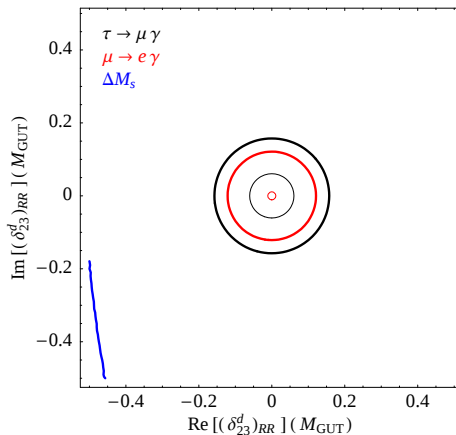


$m_0 = 210$ GeV

Gaugino/scalar mass dependence



$m_0 = 400 \text{ GeV}$



$m_0 = 210 \text{ GeV}$

Summary of Numerical Results

- For smaller scalar mass (at GUT scale)
Lepton sector constraint is stronger
- For larger scalar mass (at GUT scale)
Hadron sector constraint becomes comparable.
- In both cases, future LFV experiments will explore yet unconstrained regions of parameter space.
Discovery of LFV in charged lepton sector can be expected.