INGRID Analysis Technical Note

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Abstract

In this note we summarize the INGRID analysis results with 2010a data. We measured the neutrino event rate, the beam profile center and these stability for the confirmation and support of 2010a oscillation analysis. We select the neutrino interaction, mainly charged current interaction, at each module and reconstruct the neutrino beam profile. We compare some distributions between data and MC and found good agreement. We get the data/MC ratio for the event rate to be 1.072 ± 0.001(stat.) $\pm$ 0.040(syst.). The center of the neutrino beam profile found to be 0.2 ± 1.4(stat.) $\pm$ 9.2(syst.) cm for X profile and -6.6 $\pm$ 1.5(stat.) $\pm$ 10.4(syst.) cm for Y profile. The neutrino beam direction is measured as a direction from the beam origin to the center measured by INGRID and measured to be $-0.22 \pm .37$ mrad.

Contents

Chapter 1

Introduction

INGRID is on-axis near detector which consists of identical 14 modules [1] to monitor the beam stability. Each module has a sandwich structure made of iron plates and scintillator trackers.

We count the number of neutrino interactions, mainly CC interaction, at each module. The neutrino event rate is monitored and the profile is reconstructed. Fig.1.1 shows a typical event in an INGRID module. Detector coordinates are shown in the figure. INGRID uses a right-handed Cartesian coordinate system in which the z axis is the beam direction and the y axis is the vertical upward direction.

This article shows the measurements of

(1) neutrino event rate and its stability

(2) neutrino beam profile center and its stability

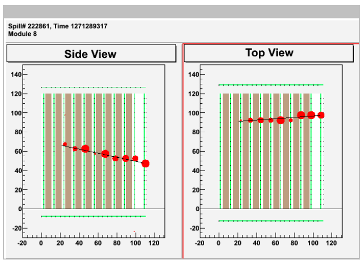

Figure 1.1: The typical neutrino event

This article is organized as follows. Chap.2 explains the overview of Monte Carlo simulation and Chap.3 describes the data set of 2010a. Chap.4 explains

[1] Additional two off-center modules and a proton module are installed after 2010a data taking.

the neutrino event selection. Finally the result of the event rate measurement and beam profile measurement are shown in Chap. 5 and Chap. 6, respectively.

Chapter 2

Monte Carlo simulation

In this chapter, we explain Monte Carlo (MC) simulation used in this analysis.

We use three MC simulation program : jnubeam, NEUT and detector simulation (Fig.2.1).

- Neutrino flux : jubeam (version 10d)
- Neutrino interaction to the target : NEUT (version 5.0.6.)
- Detector response simulation based on GEANT4 [1]

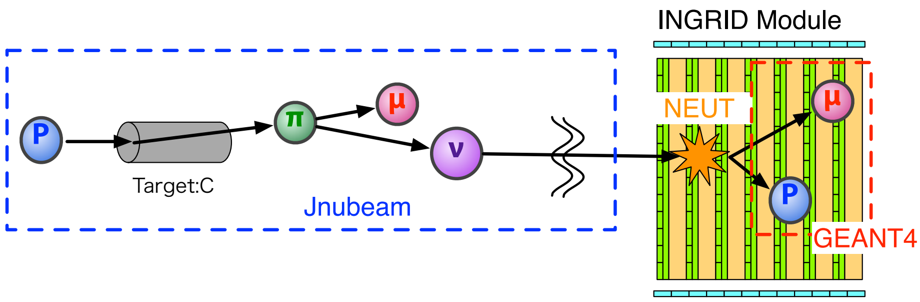

Figure 2.1: INGRID MC overview

Neutrino flux prediction

Predictions for the neutrino flux are obtained via jnubeam version 10d. For a detailed description of jnubeam, see Ref.**??**. At the INGRID detector location a total neutrino flux per proton on target of *** cm^{-2} is expected and the flux is dominated by muon neutrino (%). Figure 2.2 shows the energy spectrum at

[1]This INGRID MC is not the software of ND280 software packages

the INGRID center module, horizontal edge module and the Super-Kamiokande detector. Because each INGRID module covers the different off-axis angle, neutrino energy spectrum is slightly different at each module. Table 2.1 shows the mean neutrino energy at each module and the mean energy is different for about 0.3 GeV between the center module and the edge module.

Figure 2.2: Neutrino energy spectrum predicted by jnubeam

module#	0	1	2	3	4	5	6	SK
Mean energy[GeV]								

Table 2.1: Mean of the beam neutrino energy

Neutrino interaction simulation

We simulate neutrino interactions with target iron plate (Fe) in the INGRID detector with the NEUT program libraries. Currently, we simulate all the neutrino interaction with iron (Fe) as the target nuclei, although INGRID consists of iron, scintillator and support material. Simulation with correct material is under preparation. For a detailed description of NEUT, see Ref.**??**.

Detector response simulation

We simulate the detector response to the generated particles from the neutrino interaction with simulator based on GEANT4. We obtain the x and y position of the neutrino interaction from jubeam flux file. The vertex z is uniformly generated in iron and scintillator tacker taking into account the mass ratio of iron planes (99.54 ton) to scintillator planes (3.74 ton).

Detector response simulation includes following effects which have an impact to the efficiency to neutrino interactions.

- We tuned the conversion factor from energy deposit to number of photon in MC simulation by adjusting the peak PE of beam related sand muon. We also includes the quenching effect of scintillator, attenuation of photon propagating in the fiber and MPPC response model (we refer the report by Calibration group of ND280 working group put on `http://www.t2k.org/nd280/calib/Meetings/Jan10Workshop/MPPClinearity/at_download/filet2k.org`). Figure 2.3 shows the typical PE distribution of beam related sand muon after these MC tuning.

- Although the inefficiency resulted from the photoelectron statistics (light yield is $\sim$ 30 PE at peak and threshold is 2.5 PE) is expected to be small (0.1%), each channel has $1 \sim 2\%$ inefficiency resulted from the gap between scintillator bars, which is studied by cosmic-ray data. Figure 2.4 black line shows the tracking inefficiency of data. Because the particle with low angle can go through the gap easily, the inefficiency become larger to the low angle track. We tuned the cross-section of the scintillator bar in MC with real geometry (Fig.2.5) to reproduce the inefficiency of data (Fig.2.4 blue and red).

- Accidental hits by MPPC dark noise result in miss reconstruction of the track, miss identification of the vertex position for example. In consequence the number of neutrino event candidates should be influenced by MPPC dark noise. In MC the MPPC noise hit is generated to reproduce the number of PE, timing, and noise rate of data (Fig.2.6 shows the timing distribution).

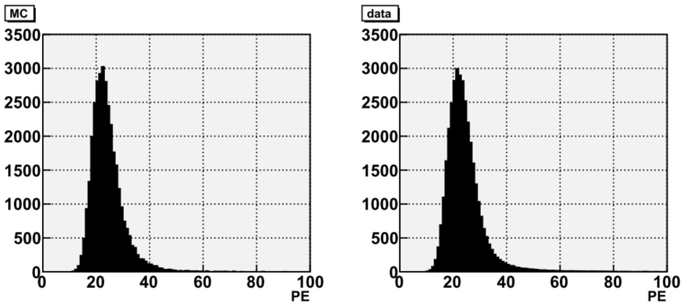

Figure 2.3: PE distribution of beam related sand muon. Left plot is MC simulation and right one is data.

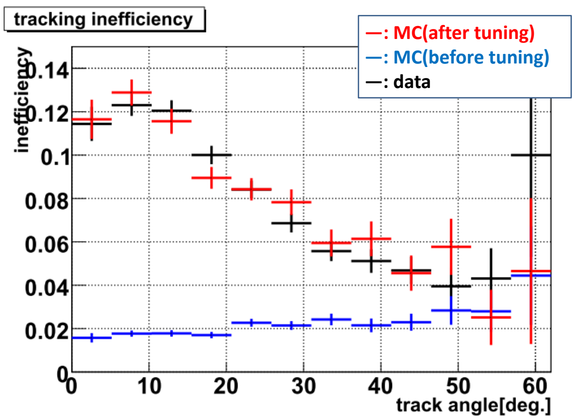

Figure 2.4: Tracking inefficiency as a function of angle of the reconstructed track

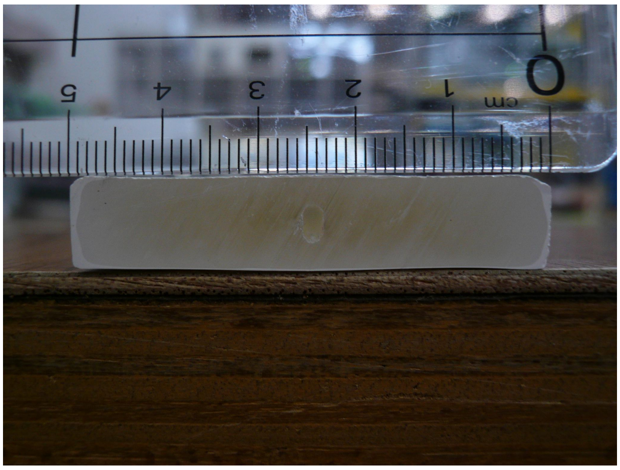

Figure 2.5: Photo of the cross-section of the scintillator bar

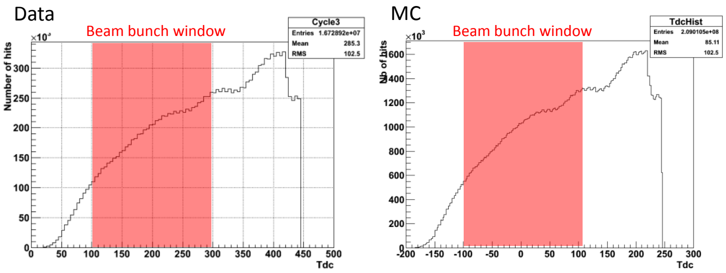

Figure 2.6: The timing distribution of MPPC noise. left is data and right is MC.

Chapter 3

Data set

We took beam data from January to June, 2010. Data taking period, number of good spills and number of INGRID good spills are summazied int Table 3.1. Data taking efficiency for entire period is 99.9%, and total number of delivered protons is 3.255×10^{19}.

The MC simulation corresponds to an equivalent of 100×10^{21} pot.

MR run #	Period	Good spills	INGRID good spills	Protons at CT05
29	Jan. 23 - Feb. 5	26813	26813	0.32×10^{18}
30	Feb. 24 - Feb. 28	59256	59070	1.12×10^{18}
31	Mar. 19 - Mar. 25	86980	86935	1.97×10^{18}
32	Arp. 14 - May. 1	237350	236647	7.64×10^{18}
33	May. 9 - Jun. 1	350079	350012	1.22×10^{19}
34	Jun. 7 - Jun. 26	246504	246410	9.30×10^{18}
Total		1006982	1005887	3.26×10^{19}

Table 3.1: Summary of data sets

Chapter 4

Neutrino event selection

4.1 Event selection

We select a long track of charged particle started within the fiducial volume of an INGRID module to select the neutrino interaction. Before reconstruction of the track, plane activity and photo-electron (PE) cut are applied to reject an accidental noise event. After reconstruction of the track, VETO cut and fiducial cut are applied to reject the incoming particle from the neutrino interaction at upstream materials. The order of event selections is shown below.

(1) Time clustering

(2) Number of active planes > 2

(3) PE/active layer > 6.5

(4) Tracking

(5) Track matching

(6) Beam timing cut

(7) Upstream VETO cut

(8) Fiducial volume cut

All selections are done a module by module and bunch by bunch basis. In this analysis, the channel which has a ADC signal larger than 2.5 PE, which corresponds to TDC threshold, is defined as the hit.

At the first step hits are clustered with following criteria; When more than 3 hits are found within 100 nsec, all the hits within $\pm$ 50 nsec from the mean time are grouped in a cluster. Within the cluster the number of planes with at least one coincidence hit in a both x and y layers [1] , which are called active

[1] INGRID module consists 11 planes and the plane consists 2 layers. Each layer has 24 scitillator bars and the direction of scitillator is perpendicular each other layer.

layers, is counted. Figure 4.1 shows the distribution of number of planes with the active layers, which are called active planes, and Fig.4.2 shows total PE of all hits in the active layer divided by the number of active layers (PE/(number of active layers)). The event with more than 2 active planes and more than 6.5 PE/(number of active layers) is selected.

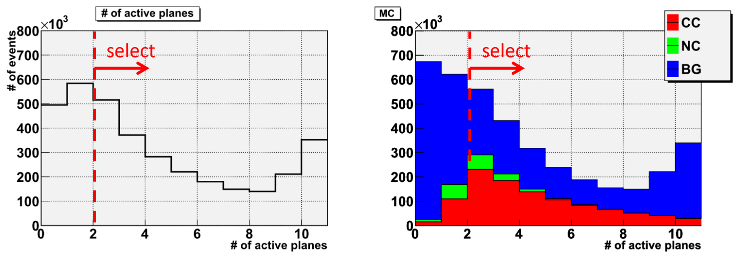

Figure 4.1: The number of active planes(left:DATA, right:MC normalized by area)

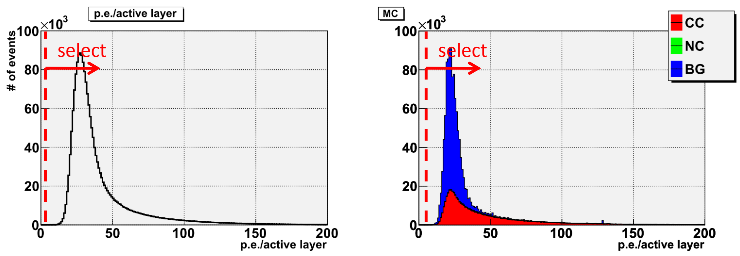

Figure 4.2: PE/(number of active layers) (left:DATA, right:MC normalized by area)

After these selections track is reconstructed. First the hits in the most downstream active plane are adopted as end-point of the track. Looking at the hits in next upstream plane in order, the hit is selected as track if calculated slope is matched with straight line.

After reconstruction of the track some badly fitted tracks are rejected by considering the difference of vertex z between a 2-D track in x view and y view. Fig.4.4 shows the distribution of the difference of the vertex z between 2-D track in X-view and Y-view. We require the difference is smaller than 2 planes. Figure 4.3 shows the distribution of the reconstructed angle.

Because there are some background events such as cosmic-ray on beam off timing, the events within $\pm$ 100 nsec from the expected timing of each bunch are selected (Fig.4.5).

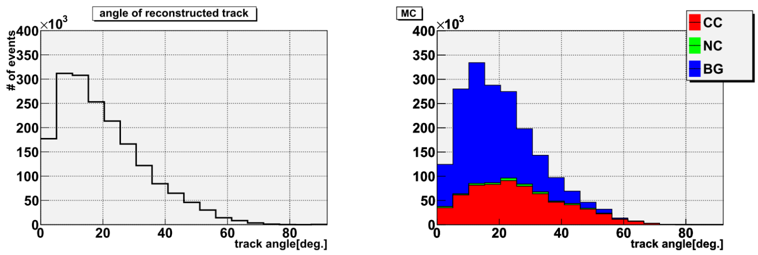

Figure 4.3: Angle of reconstructed track(left:DATA, right:MC normalized by area)

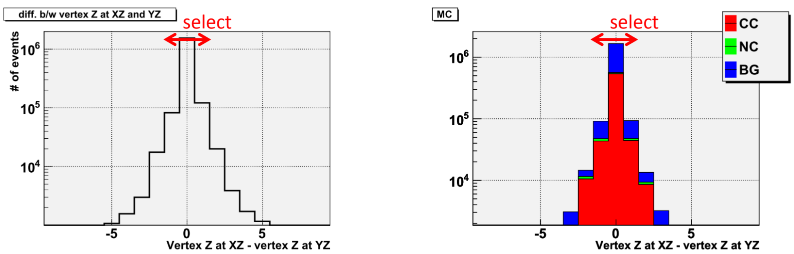

Figure 4.4: Difference of the vertex z in X-view and Y-view(left:DATA, right:MC normalized by area)

Finally we apply two selections to reject the incoming particles produced by the neutrino interactions in upstream materials. First one is upstream VETO selection. If the VETO plane has a hit at the position expected from the reconstructed track, its event is rejected. Figure 4.6 shows the example of the event rejected at this selection. After that we apply fiducial volume cut. The fiducial volume is the cubic volume which is defined as $\pm$ 50 cm from an INGRID module (Fig.4.7). The most upstream hit position associated with the reconstructed track is applied as the vertex and we require that vertex is in the fiducial volume (Fig.4.8 and 4.9).

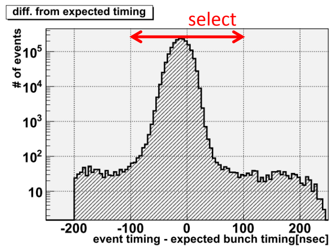

Figure 4.5: Time residual plot

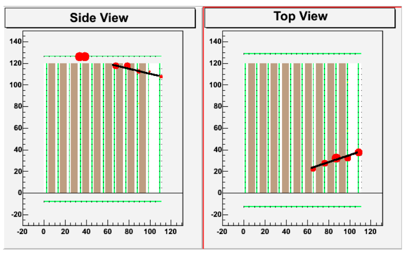

Figure 4.6: The event rejected by upstream VETO selection

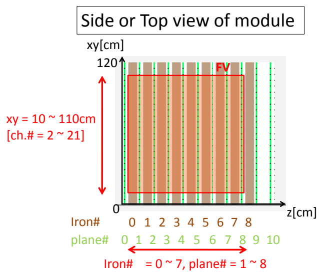

Figure 4.7: The definition of fiducial volume

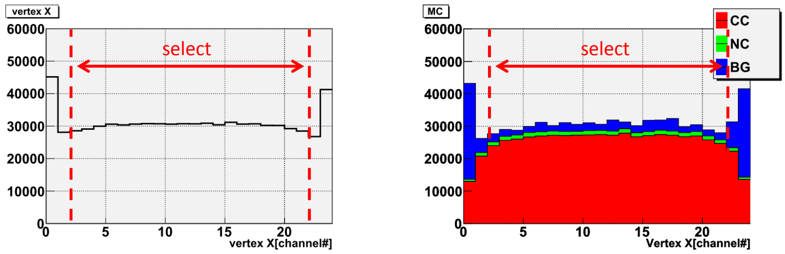

Figure 4.8: vertex x(left:DATA, right:MC normalized by area)

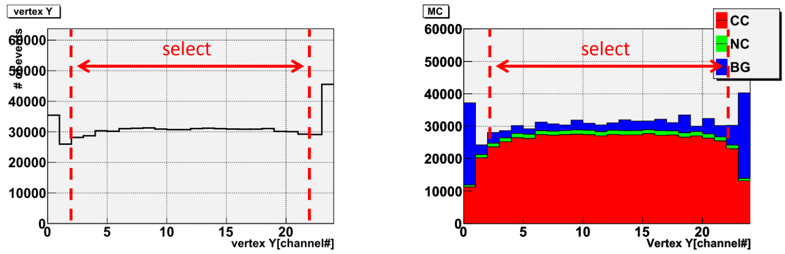

Figure 4.9: vertex y(left:DATA, right:MC normalized by area)

Event selection summary

The number of events and selection efficiencies at each selection step is summarized in Table 4.1. We obtained 493813 neutrino event candidates among 2010a data set.

	selection	Data		MC	
1	# of active planes > 2	1906146		1.97×10^6	
2	PE / active layers > 6.5	1906078	(1.00)	1.97×10^6	(1.00)
3	Tracking	1804786	(0.95)	1.83×10^6	(0.93)
4	Track matching	1749548	(0.97)	1.77×10^6	(0.97)
5	Beam timing	1747181	(0.99)	1.77×10^6	(1.00)
6	Upstream VETO cut	745912	(0.43)	7.35×10^5	(0.42)
7	Vertex in fiducial	493813	(0.66)	4.75×10^5	(0.66)

Table 4.1: Summary of the event selection. Data and MC are normalized by pot

4.2 Basic distribution of data and MC simulation

In this section we show some distributions of selected events. In each distribution, there are two plots; one is data and MC simulation overlaid and one is data/MC. In each plot the distribution of MC simulation is normalized by the area of the distribution of data. We found good agreement between DATA and MC.

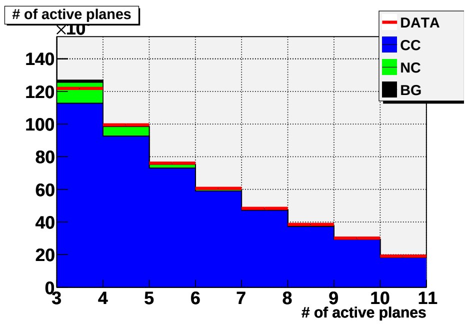

Figure 4.10: number of active planes

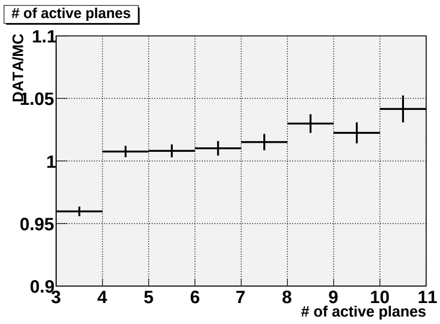

Figure 4.11: DATA/MC

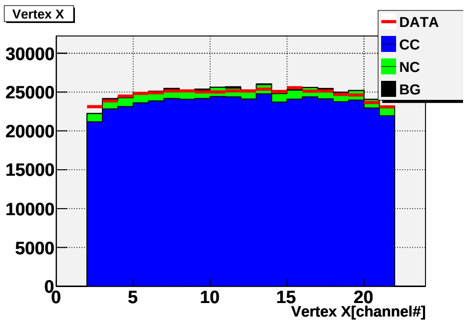

Figure 4.12: Vertex X

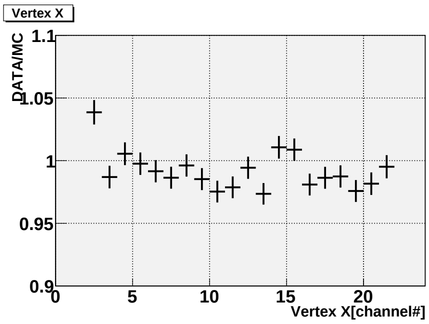

Figure 4.13: DATA/MC

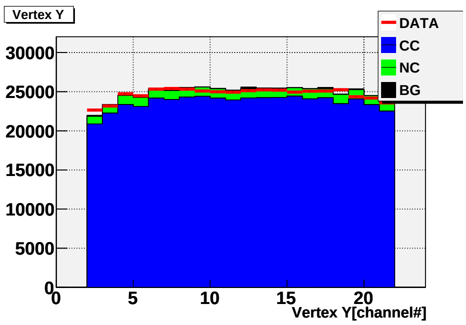

Figure 4.14: Vertex Y

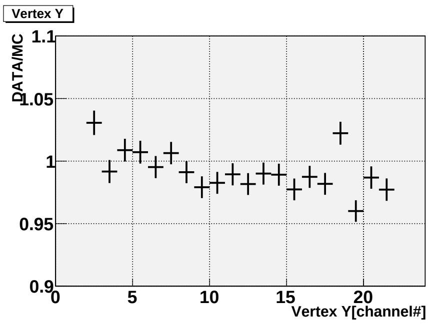

Figure 4.15: DATA/MC

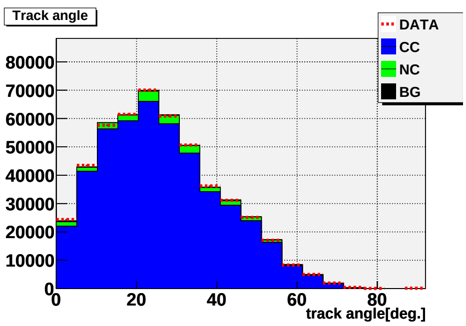

Figure 4.16: Reconstructed track angle

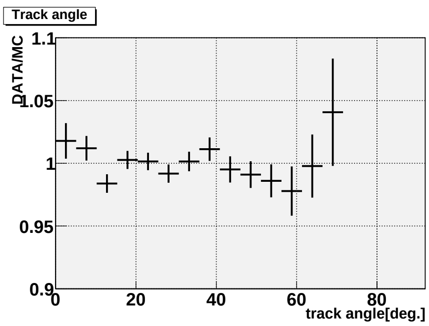

Figure 4.17: DATA/MC

4.3 Reconstruction resolution

Reconstruction resolution is checked by MC to compare the reconstructed value and the MC true information. The results of vertex X, Y and track angle show Fig.4.18 and 4.19. Their r.m.s. for CCQE events are 2.7 cm for X, 2.8 cm for Y and 3.8 degree, respectively.

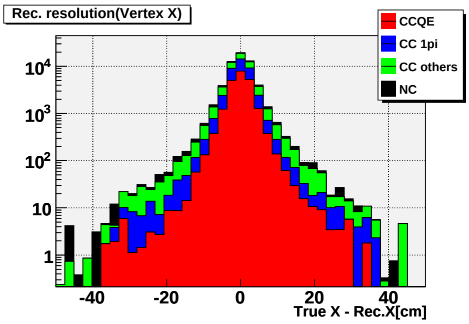

Figure 4.18: X resolution

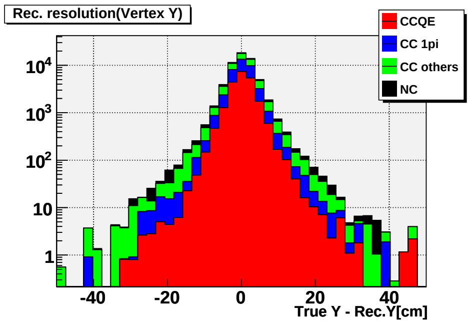

Figure 4.19: Y resolution

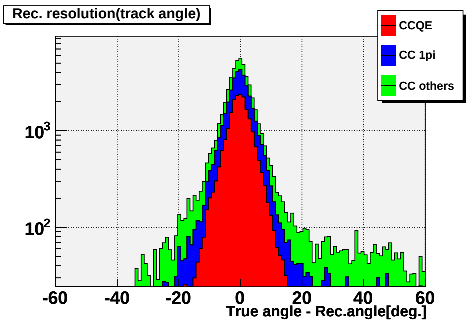

Figure 4.20: Angular resolution of reconstructed track

4.4 Efficiency to neutrino interaction

The event selection efficiency as a function of true neutrino energy is shown in Fig.4.21 and 4.22. Here, the efficiency is defined as the ratio of the number of selected events to that of the events generated inside the fiducial volume. Table 4.2 shows the selection efficiency for each module. Because the energy spectrum of the beam neutrino is slightly different module by module, the selection efficiency is also different.

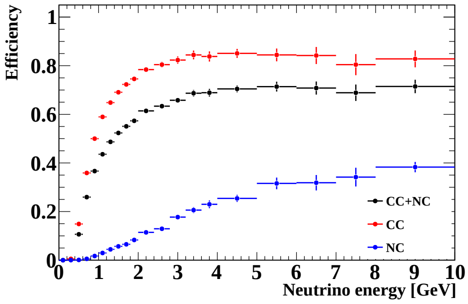

Figure 4.21: Neutrino event selection efficiency

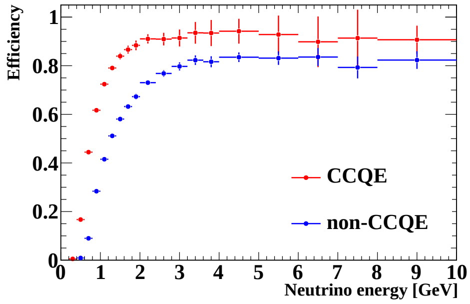

Figure 4.22: Selection efficiency for CCQE and CC others.

module	Mean efficiency[%]
0	51.7
1	54.0
2	55.1
3	55.1
4	55.0
5	54.2
6	51.2
7	52.6
8	54.4
9	55.1
10	55.0
11	54.6
12	54.1
13	51.8

Table 4.2: Mean efficiency of each module

Chapter 5

Event rate measurement

5.1 Event rate stability

Fig.5.1 shows daily event rate normalized by delivered pot. We succeeded to measure the daily event rate with about 1.7% statistical error each day. The chi-square calculated from the average rate is 86 for 76 degrees of freedom. It is concluded that the beam events is stable within statistical error.

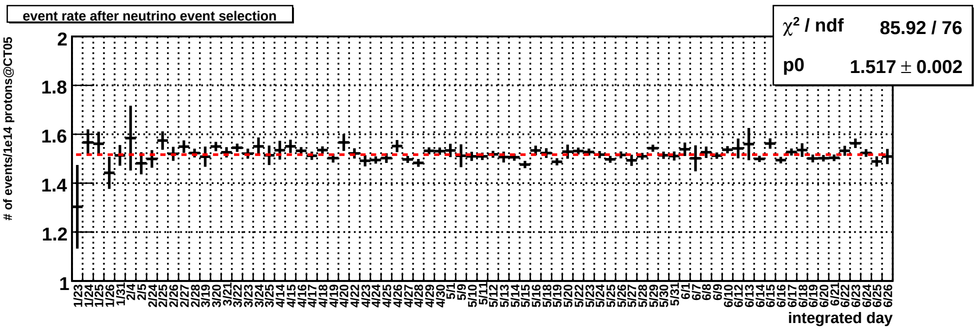

Figure 5.1: The stability of daily event rate

5.2 data/MC of event rate

To obtain the number of events in the fiducial volume ($N^{obs.}$), we need to do following corrections.

(1) Accidental MPPC noise

(2) Iron mass

(3) Beam related background

Detail description is in Chap. A.

We derive the formula to evaluate the number of events in the fiducial volume from number of selected events ($N^{sel.}$); $N^{obs.} = N^{sel.} \times \frac{1}{1+C}$, where C is the correction factor. The corrected number of events are summarized in Table 5.1 respectively. Detail for it is summarized in Chap. A. Finally we obtain DATA/MC to be 1.072 ± 0.001 (stat.).

Number of selected events	493813
Corrected number of events	508511

Table 5.1: Number of events before and after corrections

5.3 Systematic error of event rate

Table 5.2 shows the systematic errors which does not include physics uncertainty such as neutrino cross-section.

Item	Error[%]
Accidental MPPC noise	0.7
Iron mass	0.1
Beam related background	0.2
Fiducial selection	1.1
Hit efficiency	1.8
Tracking efficiency	1.4
Track matching selection	2.7
Not beam-related background	<0.1
p.e./active layer selection	<0.1
Beam timing selection	<0.1
Total	3.7

Table 5.2: Systematic error table

Accidental MPPC noise

The effect of MPPC noise is studied with MC in which the MPPC noise hit is generated to reproduce number of PE, timing, and noise rate of DATA. We found that the more MPPC noise rate is, the more neutrino events are lost due to miss identification of vertex Z or miss counting of number of active planes. Its effect is found to be linear and slope is estimated to be -0.9585 %. Two sources of systematic errors are considered. First one comes from the error on the linear fit. To get this systematic error, we multiply the fit error by the maximal measured noise rate. Second one comes from the measurement of noise. Correction factors are calculated using the average noise rate measured on one period. But this noise rate fluctuates in time (probably due to temperature variations). So we measure the maximal difference between average noise rate and noise rate measured at different times during one period, and using the linear relation between noise rate and variation of number of events we get the systematic error. The quadratic sum of these two errors is 0.7 %.

Iron mass

Before construction of INGRID the mass of each iron plate was measured with a precision of 1 kg, which corresponds to 0.13 % of the mass of one iron plate. We will use this figure as the systematic error on this correction factor. We might need to increase this systematic error in the future, as the correction factors are calculated using the mass of the whole iron plate, when we actually use a fiducial cut in analysis, only interactions in the central part of the iron plates are kept.

Beam-related background

We estimated the contamination fraction of beam related background with wall neutrino Monte Carlo. The fraction is estimated to be 0.4% , in which the number of interactions of background is normalized to compare the number of dirt muon in DATA and MC. There is a 35% difference from POT expectation, which is considered as one of the source of the systematic error. We considered 20% neutrino flux uncertainty and 20% cross section uncertainty as other sources of the systematic error. Finally 0.2% ($=\sqrt{0.35^2 + 0.2^2 + 0.2^2}$) is applied as the systematic error.

Fiducial selection

To estimate the uncertainty of fiducial selection and the effect of non uniformity of iron plate, we divided fiducial in several horizontal slices and checked the difference between DATA and MC. Table 5.3 shows the result. The maximum absolute value, 1.1%, is applied as systematic error.

selection	DATA	MC	|DATA - MC|
<50 cm from center(nominal selection)	100.0	100.0	0.0
<25 cm	25.6	25.2	0.4
$25 \sim 40$ cm	39.9	39.3	0.6
$40 \sim 50$ cm	34.4	35.5	1.1
Systematic error (Maximum absolute)			1.1

Table 5.3: DATA-MC for several sub fiducial volume

Hit efficiency

We estimated the relation between hit efficiency and number of selected events with MC. Fig. 5.2 shows the result from which the systematic error of hit efficiency is estimated to be 1.8% because hit efficiency has 1.1% uncertainty. [1]

Track matching selection

In the neutrino event selection, after reconstruction of XZ track and YZ track we require track start point matching. To estimate the uncertainty of the selection, we changed the tolerance for the matching and checked the difference of the number of selected events between DATA and MC. Table 5.4 shows the result. The maximum absolute value, 2.7%, is applied as systematic error.

[1] 0.5% of the measurement error of hit efficiency, 1.0% of the tuning of hit efficiency in MC

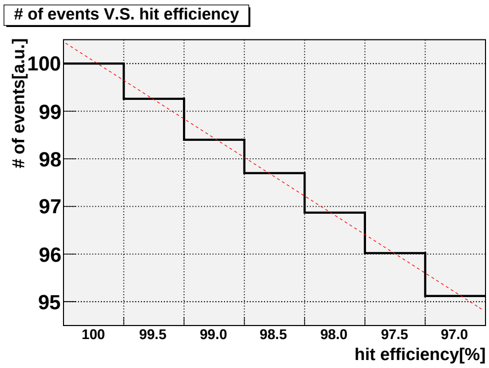

Figure 5.2: hit efficiency V.S. number of selected events

Vertex Z of XZ track - Vertex Z of YZ track	DATA	MC	|DATA - MC|
-1, 0, +1 (nominal selection)	100.0	100.0	0.0
0	83.0	85.7	2.7
-2, -1, 0, +1, +2	104.0	103.0	1.0
Systematic error (Maximum)			2.7

Table 5.4: DATA-MC for several tolerance of track matching.

Tracking efficiency

To check the difference of the tracking efficiency between DATA and MC, the tracking efficiency is compared with several sub-sample selected by number of active planes. Table 5.5 shows the result. The maximum absolute value, 1.4%, is applied as systematic error.

Not beam-related background

The off-bunch data (cycle 17 $\sim$ 22 where as on-bunch cycle is 4 $\sim$ 9) are analyzed with same procedure and only 93 events are selected whereas the number of signal is 493813. It is negligible.

PE/active layer selection

To estimate the uncertainty of PE/active layer selection, we changed the cut value and checked the difference of number of selected events from one with nominal cut. The result is the difference is less than 0.01% and its uncertainty is negligible.

number of active planes	DATA	MC	|DATA - MC|
3	87.6	86.9	0.7
4	93.2	91.8	1.4
5	94.7	94.3	0.5
6	95.6	96.2	0.6
7	96.2	96.6	0.4
8	96.7	96.8	0.1
9	98.7	97.9	0.8
10	99.1	99.0	0.1
Systematic error (Maximum)			1.4

Table 5.5: The tracking efficiency of DATA and MC with several sub sample

beam timing selection

To estimate the uncertainty from neutrino beam timing, we changed the cut value and checked the difference of number of events from nominal cut. The difference is less than 0.01% and it is negligible.

Chapter 6

Measurement of beam profile

6.1 Stability of beam profile

Fig.6.1 shows horizontal and vertical beam profile with RUN 32 data. Fitted center with gaussian is 0.1 ± 2.9 cm for horizontal and -10.9 ± 3.2 cm for vertical. Fig.6.4 and 6.5 show the monthly beam center of horizontal and vertical respectively. We succeeded to measure the profile center with about 4.2 cm statistical error for each month. The chi-square is calculated to be almost one to the degree of freedom (0.8 for X and Y center). It is concluded that the beam profile center is stable within statistical error.

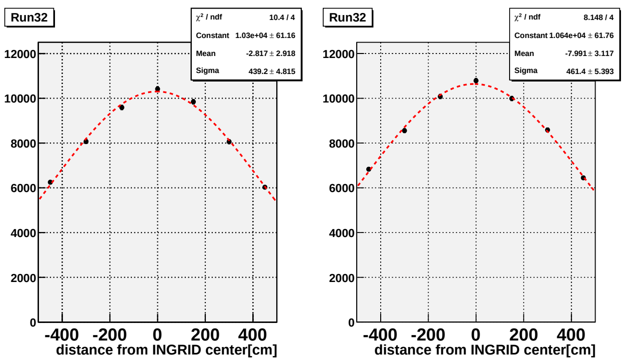

Figure 6.1: Horizontal profile(left) and vertical profile(right)

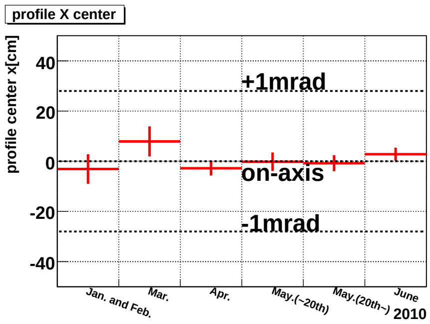

Figure 6.2: Horizontal profile center.

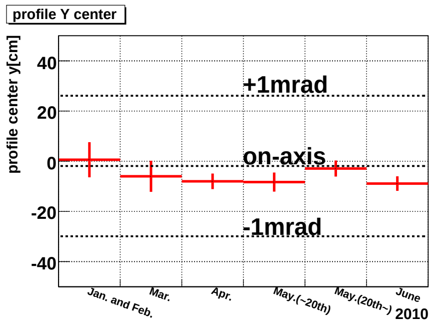

Figure 6.3: Vertical profile center

6.2 The systematic error of beam center

We estimated the systematic error with toy profile MC in which the number of events at each module is varied with 3.7% from original profile made by RUN 29 34 all data. 100'000 profiles are generated and the RMS of fitted center is applied as systematic error. The result shows 9.2 cm (0.33 mrad) for horizontal center and 10.4 cm (0.37 mrad) for vertical center.

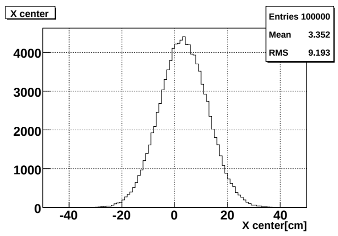

Figure 6.4: Fitted Horizontal center with 100'000 profiles

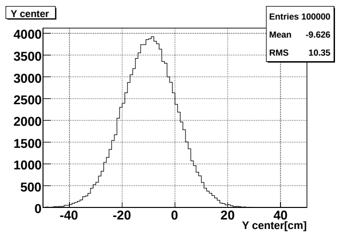

Figure 6.5: Fitted Vertical center with 100'000 profiles

Chapter 7

Conclusion

In this note we have presented the measurement of the neutrino event rate and profile center in INGRID durint 2010a. We selected the neutrino interactions to reconstruct the long track started within fiducial volume. The results have been compared to MC and found good agreement with DATA. Finally DATA/MC of the event rate and beam profile centers have been evaluated with an associated systematic error:

$$\begin{aligned} \mathrm{R}_{\mathrm{DATA/MC}} &= 1.072 \pm 0.001(\mathrm{stat.}) \pm 0.040(\mathrm{syst.}) & \\ \mathrm{Xcenter} &= +0.2 \pm 1.4(\mathrm{stat}) \pm 9.2(\mathrm{syst.}) & \mathrm{cm} \\ \mathrm{Ycenter} &= -6.6 \pm 1.5(\mathrm{stat.}) \pm 10.4(\mathrm{syst}) & \mathrm{cm} \end{aligned}$$

Appendix A

Correction factors for neutrino event rate

Iron mass

In INGRID most of the neutrino interactions occur in the 9 iron targets of each module. During their fabrication, there was a tolerance on the thickness of those iron planes. This results in iron planes having slightly different volumes, and as a consequence different masses. The maximal variation from design mass is 2.15 % from the given tolerance on thickness. The mass of each iron plane was measured at the end of the fabrication process, so we can deduce correction factors for the expected number of events for each module, by using the fact that 95.2 % of interactions in one module occur in the iron.

Accidental MPPC noise

Another correction on the number of observed events comes from noise hits in the detector. Those noise hits reduce the number of reconstructed events compared to the case when there is no noise. To correct this effect, we use the following procedure:

(1) Measure noise in data

(2) Create a noise simulation to reproduce those measurements

(3) Use Monte Carlo simulation to compare the number of reconstructed events with and without adding noise

(4) Deduce from the simulation correction factors and systematic errors

Noise is measured in beam data. We measure the rate of noise hits, which are defined as hits occuring in the detector when no particles are actually going through the detector. To find such hits, we look at cycles where beam spills are not coming (INGRID records data on 23 integration cycles, but beam spills only arrive during 6 of them), and perform regular event selection to make sure there is no cosmic particle in the detector. We then measure a noise rate for each channel of the detector, as well as light yield and timing distribution.

Noise is then simulated with a given probability for each channel. Timing for the noise hits is simulated using the distribution measured in data. Light yield is then simulated using a measured light yield distribution for the corresponding timing.

Monte Carlo simulation is then used to measure the variation of number of reconstructed events due to noise. We first reconstruct events on Monte Carlo files which do not include noise hits, then add noise hits to those files and perform reconstruction again. We then compare the number of reconstructed events in each case. The simulation is using jnubeam 10c.

From this simulation we have for each module a noise rate and the variation of number of reconstructed events due to noise. There is a linear relation between them as can be seen on Fig.A.1. We will use this linear relation to make corrections on the number of observed events. This relation is:

$$\text{Variation of number of events [\%]} = -0.9585 * \langle \text{noise rate} \rangle$$

Those corrections are made for each module, and each subset of events we are considering. In each case we measure the noise rate, and then from the linear relation deduce the variation of number of reconstructed events which should be used as a correction factor.

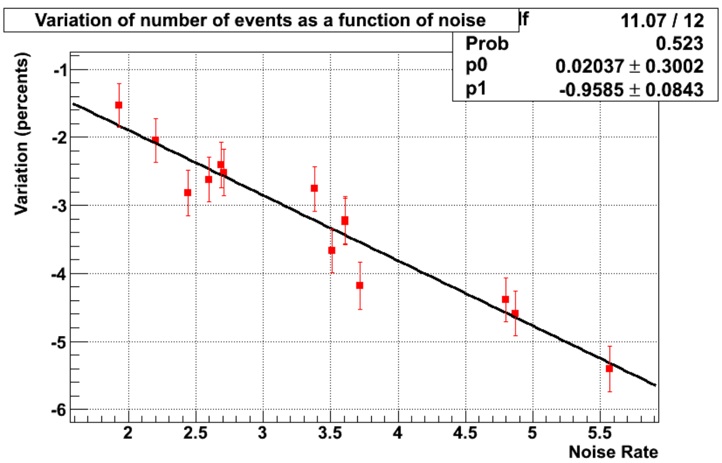

Figure A.1: Variation of number of reconstructed events as a function of noise rate

Beam related background

We estimated the contamination fraction of beam related background with background MC in which neutrino and the interaction is generated in upstream dirt (10 x 10 x 5 m^3).

Almost all contaminations come from short track induced neutron ($\sim 50\%$) or gamma ($\sim 40\%$) and dirt muon ($\sim 10\%$) which is not detected accidentally due to scintillator inefficiency.

In background MC, number of generated interactions is normalized so that number of rejected events at upstream VETO selection, which consists dirt muon mainly, is equal to DATA and MC. The difference from POT expectation is 35% and it is considered as systematic error. Finally contamination fraction is estimated to be 0.4% and it is applied as one of the correction factor.

Summary of the correction factor

Run by run and module by module correction factors are summarized int Table A.1

	module	29,30	31	32	33	33	34
number of events	0	1054	1548	5956	4119	5425	6962
	1	1526	2033	7827	5432	7122	9520
	2	1875	2476	9360	6492	8555	11622
	3	1882	2570	10133	6795	9078	12191
	4	1831	2459	9627	6636	8683	11651
	5	1524	2176	7876	5421	7217	9588
	6	1058	1585	5837	4172	5421	7285
	7	1229	1717	6636	4509	5826	8100
	8	1588	2187	8351	5819	7620	10270
	9	1884	2562	9770	6632	8766	11946
	10	1949	2681	10305	6987	9373	12473
	11	1908	2520	9771	6713	8897	11871
	12	1561	2133	8146	5512	7193	9822
	13	1218	1659	6263	4327	5815	7734
correction factor	0	-3.3	-3.3	-4.3	-4.0	-3.9	-3.9
	1	-2.6	-2.6	-2.6	-2.4	-2.4	-2.4
	2	-2.0	-2.0	-2.0	-1.7	-1.7	-1.7
	3	-2.3	-2.3	-2.3	-2.0	-1.9	-1.9
	4	-1.8	-1.8	-1.8	-1.6	-1.5	-1.5
	5	-1.9	-1.9	-1.9	-1.6	-1.5	-1.4
	6	-2.3	-2.3	-2.8	-2.5	-2.4	-2.3
	7	-2.7	-2.7	-2.5	-3.5	-3.3	-3.1
	8	-2.2	-2.2	-2.0	-3.0	-2.8	-2.6
	9	-2.1	-2.1	-2.7	-4.1	-3.8	-3.6
	10	-4.2	-4.2	-4.1	-5.4	-5.2	-5.0
	11	-1.9	-1.9	-1.8	-2.9	-2.7	-2.6
	12	-4.9	-4.9	-4.7	-6.2	-6.0	-5.9
	13	-2.5	-2.5	-2.5	-3.4	-3.2	-3.0
Number of cor.	0	1090	1601	6224	4292	5647	7242
	1	1567	2087	8037	5565	7298	9756
	2	1913	2526	9551	6606	8703	11822
	3	1927	2632	10374	6934	9255	12422
	4	1865	2504	9804	6741	8817	11825
	5	1553	2218	8028	5509	7328	9697
	6	1083	1622	6002	4278	5552	7453
	7	1263	1765	6803	4674	6024	8360
	8	1624	2237	8520	6001	7838	10545
	9	1925	2617	10040	6912	9110	12391
	10	2035	2800	10742	7385	9883	13130
	11	1945	2569	9951	6914	9143	12183
	12	1641	2242	8550	5878	7653	10435
	13	1250	1702	6421	4478	6005	7976

Table A.1: Correction factors

Bibliography

[1] will be put the name of jnubeam tech. note

[2] will be put the name of NEUT doc.